

Electromagnetic effect due to axion dynamics in a superlattice of topological insulators

Katsuhisa Taguchi (Nagoya University)

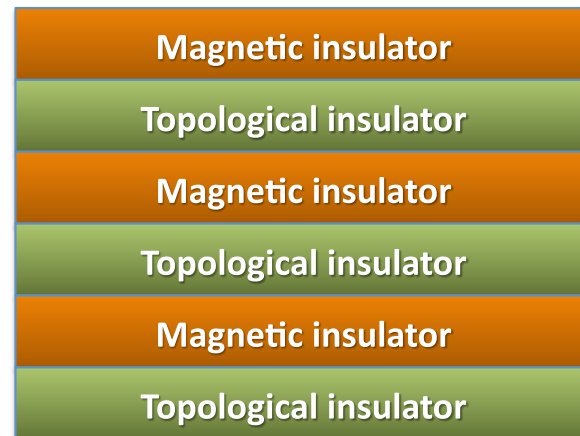


Collaborators:

Tatsushi Imaeda, Tetsuya Hajiri,
Tatsuhisa Naka, Takuya Shiraishi,
Naoya Kitajima (Nagoya University)

prepare to submit ...

Superlattice

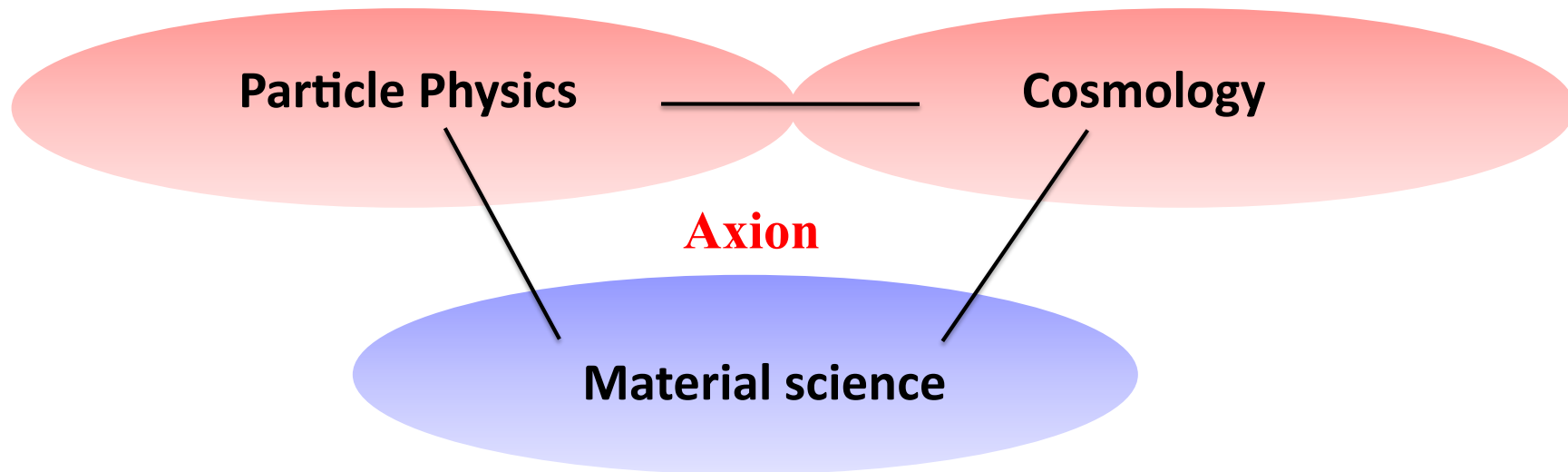


Outline

1. Introduction to **“axion”** in material science
2. Our idea to generate the dynamical “axion”
by an applying magnetic field
3. Summary

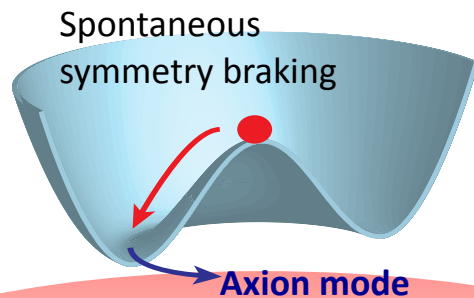
“Axion” in material science

Axion is a common research theme
in the particle physics, cosmology, and material science



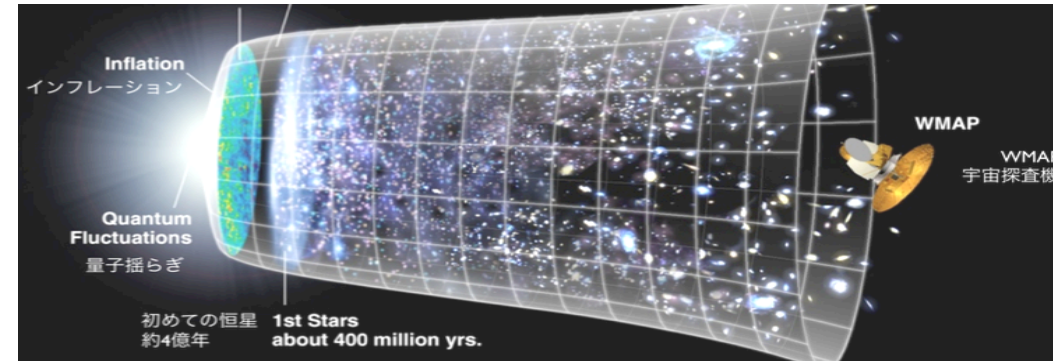
“Axion” in material science

Hypothetical particle
to resolve “strong CP problem”



Particle Physics

Candidate for Dark Matter



Cosmology

Dark Energy
(68%)

Dark Matter
(27%)

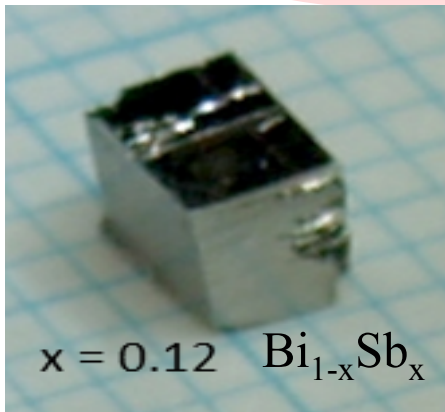
Ordinary matter
(5%)

Planck 2013 results

Axion

Material science

“Axion” [axion-like particle]
is hosted in topological materials



Topological materials

It's very hot theme for material science

Theory (~ 2000) \Rightarrow discovery of theory about "topology" in the condensed matter

2007 : Science 138(2007) 766 [**2D topological insulator**]

2008 : Nature 452(2008) 970 [**3D topological insulator**]

2016 Noble Prize in Physics



Mahmoud
David J. Thouless



Mahmoud
F. Duncan M. Haldane



J. Michael
Kosterlitz

The Nobel Prize in Physics 2016 was awarded with one half to David J. Thouless, and the other half to F. Duncan M. Haldane and J. Michael Kosterlitz *"for theoretical discoveries of topological phase transitions and topological phases of matter"*.

Topological insulators

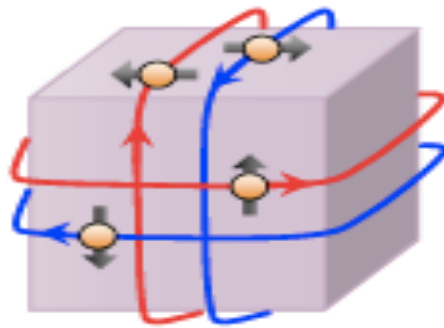
Topological insulator

- Bulk: Insulator
- Surface: Metallic

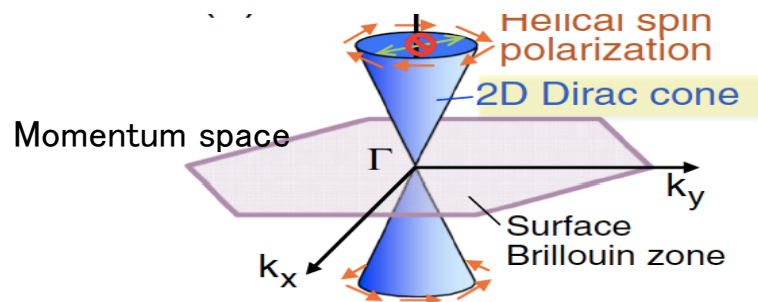
Hasan and Kane, *Rev. Mod. Phys* (2010),
Qi and Zhang, *Rev. Mod. Phys* (2011),
Ando *JPSJ* (2013)

Helical edge state

- Spin-dependent transport
- 2D Dirac cone



Real space

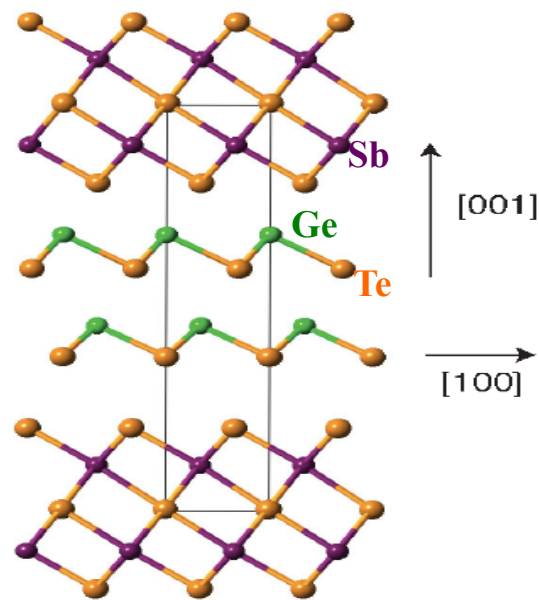


Superlattice

Superlattice of topological insulator (TI) and normal insulator (NI)

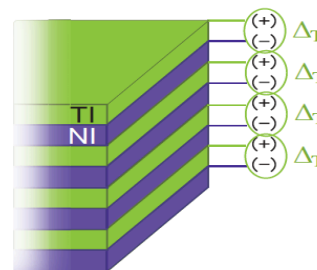
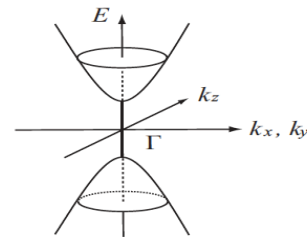
Superlattice of GeTe (NI) / SbTe (TI)

Ref: J. Tominaga, et al., Adv. Mater. Interfaces 1, 1300027(2014)



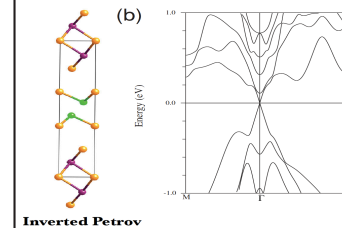
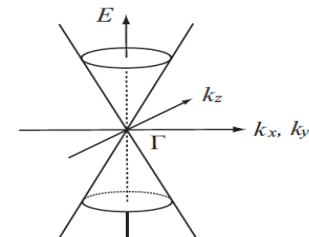
(a) $\Delta_T > \Delta_N$

3D NI
(gapped)



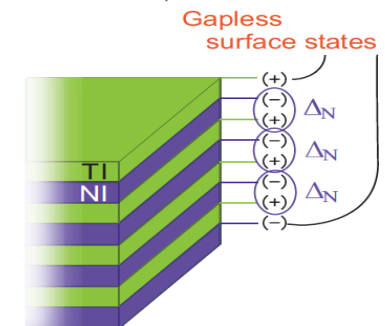
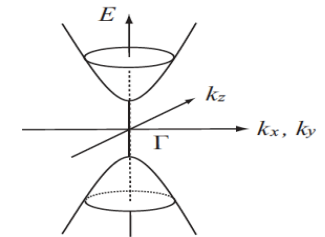
(b) $\Delta_T = \Delta_N$

Dirac semimetal
(gapless)



(c) $\Delta_T < \Delta_N$

3D TI
(bulk:gapped
surface:gapless)



This material has been already applied in *interfacial phase change memory*

Lagrangian in a topological insulator

Effective Lagrangian in a topological materials is described by

(Topological insulator, Weyl semimetal, Dirac semimetal, etc...)

Ref: F. Wilczek, PRL 58, 1799 (1987).

X.-L. Qi, T. L. Hughes, and S.-C. Zhang, PRB 78, 195424 (2008).

A. M. Essin, J. E. Moore, and D. Vanderbilt, PRL 103, 259902 (2009).

$$\mathcal{L}_a = g_a \gamma \gamma \alpha \theta(x, t) \underline{E \cdot B}$$

Diagram illustrating the components of the Lagrangian \mathcal{L}_a :

- g_a : Coupling constant
- $\gamma \gamma$: Fine structure = 1/137
- $\alpha \theta(x, t)$: \sim “axion-like” particle (or field)
- $\underline{E \cdot B}$: Electric field and magnetic field

This Lagrangian is similar to the Lagrangian in Field theory (QCD-axion)

SO

“ θ ”-term is called as “axion” field

Maxwell equations

Maxwell equation from Euler-Lagrange Equation

$$\begin{aligned}\nabla E &= 4\pi\left[\rho + \frac{e^2}{2\pi h} \nabla \left(\frac{\theta}{\pi}\right) B\right] \\ \text{rot} B - \frac{1}{c} \frac{\partial E}{\partial t} &= 4\pi\left[j + \underbrace{\frac{e^2}{2h} \nabla \left(\frac{\theta}{\pi}\right) \times E}_{j_a : \text{axion induce current}} + \underbrace{\frac{e^2}{2\pi h} \frac{\dot{\theta}}{\pi} B}_{\text{Chiral magnetic effect term}}\right]\end{aligned}$$

Anomalous Hall effect term (Hall effect w/o B)

Chiral magnetic effect term

$\nabla \theta$ and $\partial_t \theta$ couple with electromagnetic fields

Trigger characteristic transport:

Ref: F. Wilczek, PRL 58, 1799 (1987).

X.-L. Qi, T. L. Hughes, and S.-C. Zhang, PRB 78, 195424 (2008).

A. M. Essin, J. E. Moore, and D. Vanderbilt, PRL 103, 259902 (2009).

Today's topic

Study of “Axion” in **material science**

Why ?

The reasons are ...

- 1. Interaction** between “Axion” and electromagnetic fields in materials science could be large (compared with that in QCD-axion)
- 2. Controllable** “Axion” by using knowledge of material science
3. There are some analogy between particle physics (cosmology) and material science
(e.g., majorana fermion, axion ..)

Material science v.s. particle physics & cosmology

particle physics & cosmology

High energy physics

Strong CP
Dark matter

Pure research

energy
scale

context

Pure or
applied

Material science

Low energy physics
 $0.1 \text{ meV} \sim 1 \text{ eV}$

**Unconventional
electromagnetic effects**

Pure
& **Device application**
via the electromagnetic effects

Material science v.s. particle physics & cosmology

Property of “Axion” in the particle physics vs that in the material science

$$\mathcal{L}_a = g_{a\gamma\gamma} \alpha \theta(\underline{x}, t) \underline{E} \cdot \underline{B}$$

	Particle physics + cosmology	Material science
Type of axion	QCD axion, ALPs	Topological axion
Interaction strength	Very small ($\sim 10^{-19}$)	Large (~ 1)
Character	Dynamical	Not dynamical
Detection	Halo scope (cavity) + Primakoff effect	In this topic

Dynamical Axion signals in the both field have not seen yet!!

First, let's try to consider about the axion response in material science.

Take-home message

Topological materials are ...

good experimental platforms for testing axion physics,

because of

Controllable parameters such as coupling constant and
axion mass

*“Axion” is controllable in lab...
We can challenge a low-energy physics*

so

axion is **not** “dark” in material science!

Naively question:

Is there true of the effective Lagrangian even in material science ?

$$\mathcal{L}_a = g_{a\gamma\gamma} \alpha \theta(x, t) \underline{E \cdot B}$$

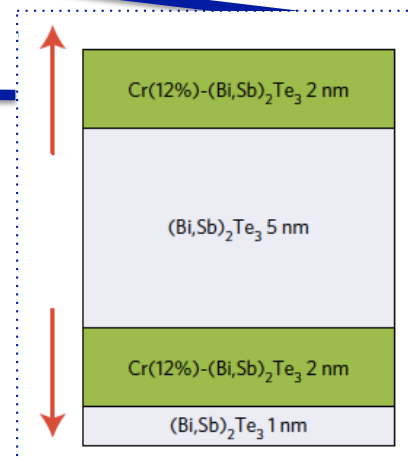
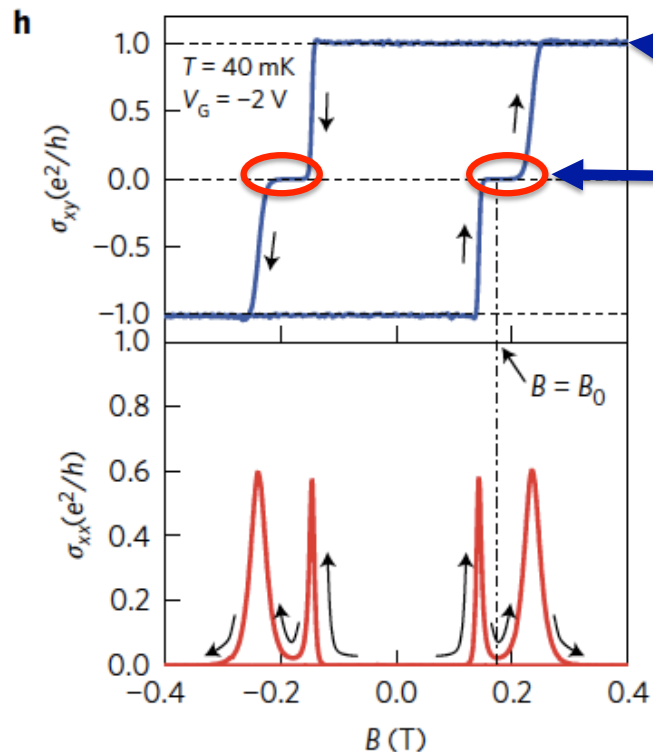
Diagram illustrating the components of the effective Lagrangian \mathcal{L}_a :

- $g_{a\gamma\gamma}$: Coupling constant
- α : Fine structure = $1/137$
- $\theta(x, t)$: \sim "axion-like" particle (or field)
- $\underline{E \cdot B}$: Electric field and magnetic field

Evidence of static axion in material science

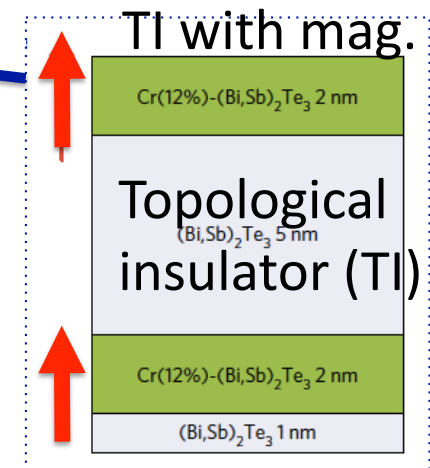
Electrical evidence: Quantum Hall effect

M. Mogi, et al., Nat Mater advance on, (2017).



Anti-parallel
magnetic configuration

There is cancellation ($=1/2 - 1/2$) of each half-quantum Hall effect on top and bottom surface of TIs



Parallel
magnetic configuration

There is no-cancellation ($=1/2 + 1/2$) half-quantum Hall effect

Optical evidence in TIs : Giant (half-quantized) Kerr effect

“Quantized Faraday and Kerr rotation and axion electrodynamics of a 3D topological insulator”

Liang Wu et al.,

Maxwell equation include axion field

Maxwell equation from Euler-Lagrange Equation

$$\begin{aligned}\nabla E &= 4\pi\left[\rho + \frac{e^2}{2\pi h} \nabla\left(\frac{\theta}{\pi}\right) B\right] \\ \text{rot} B - \frac{1}{c} \frac{\partial E}{\partial t} &= 4\pi\left[j + \underbrace{\frac{e^2}{2h} \nabla\left(\frac{\theta}{\pi}\right) \times E}_{j_a : \text{axion induce current}} + \underbrace{\frac{e^2}{2\pi h} \frac{\dot{\theta}}{\pi} B}_{\text{Chiral magnetic effect}}\right]\end{aligned}$$

However...

In usual TI, as θ is static, it does not have time-dependence. (and $\theta = \pi$)

You can not see the electromagnetic effect following Maxwell equation!!



We dare to think the dynamical θ from the analogy of particle physics

However, There is the source term generated by the time-dependent axion field

➡ How can we make such dynamical axion field ?

Our idea

“Generating the time-dependent θ using the spin-flip “

Our motivation and purpose

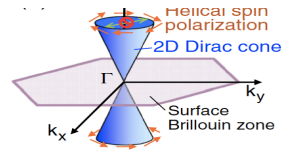
Conventional works consider **static “axion”-response**

Our work shows that ...

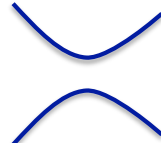
[1] how to drive **dynamical “axion”-response**

[2] the electromagnetic effect via **dynamical “axion”-response**

θ term is mathematically described when there is gap



Gapless (massless)



Gap

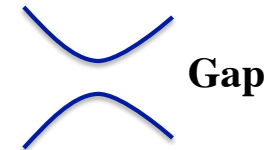
X.-L. Qi, T. L. Hughes, and S.-C. Zhang, PRB 78, 195424 (2008). [See Eq. 1]
 P. Hosur, S. Ryu, and A. Vishwanath, PRB 81, 45120 (2010).
 R. Li, J. Wang, X.-L. Qi, and S.-C. Zhang, Nat Phys 6, 284 (2010).

θ term is obtained from Hamiltonian with PT-symmetry

$$H(\mathbf{k}) = \mathbf{k} \cdot \boldsymbol{\Gamma} + \phi_4 \Gamma_4 + \phi_5 \Gamma_5,$$

$$\{\Gamma_i, \Gamma_j\} = 2\delta_{ij}$$

ϕ_4 corresponding the band gap in $\phi_5=0$



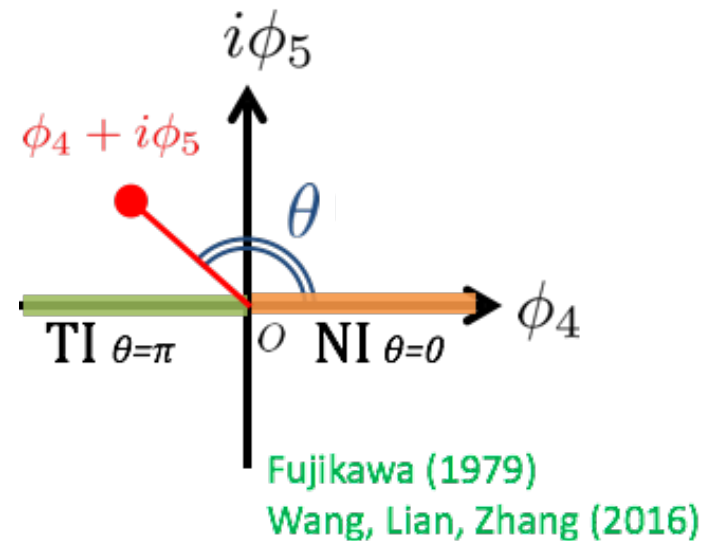
ϕ_5 corresponding the breaking P and T simultaneously

Hamiltonian with PT-symmetry

$$H(\mathbf{k}) = \mathbf{k} \cdot \mathbf{\Gamma} + \phi_4 \Gamma_4 + \phi_5 \Gamma_5,$$

Theta term

$$\theta = \text{Arctan}(\phi_5 / \phi_4)$$



[Review] P. Hosur, S. Ryu, and A. Vishwanath, PRB 81, 45120 (2010).

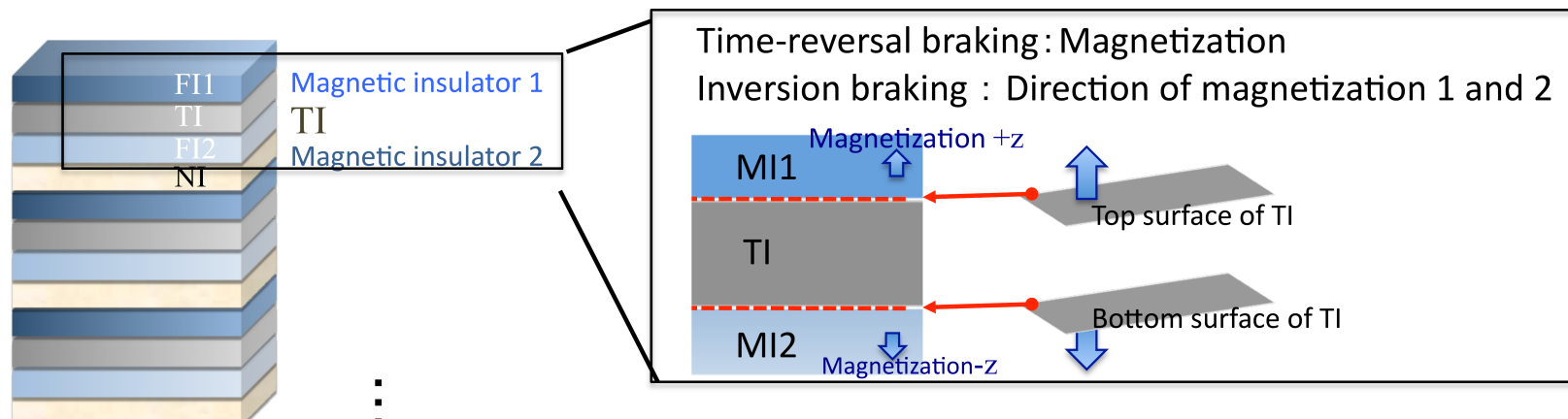
K. Fujikawa and H. Suzuki, "Path Integrals and Quantum Anomalies" Clarendon Press, Oxford, 2004

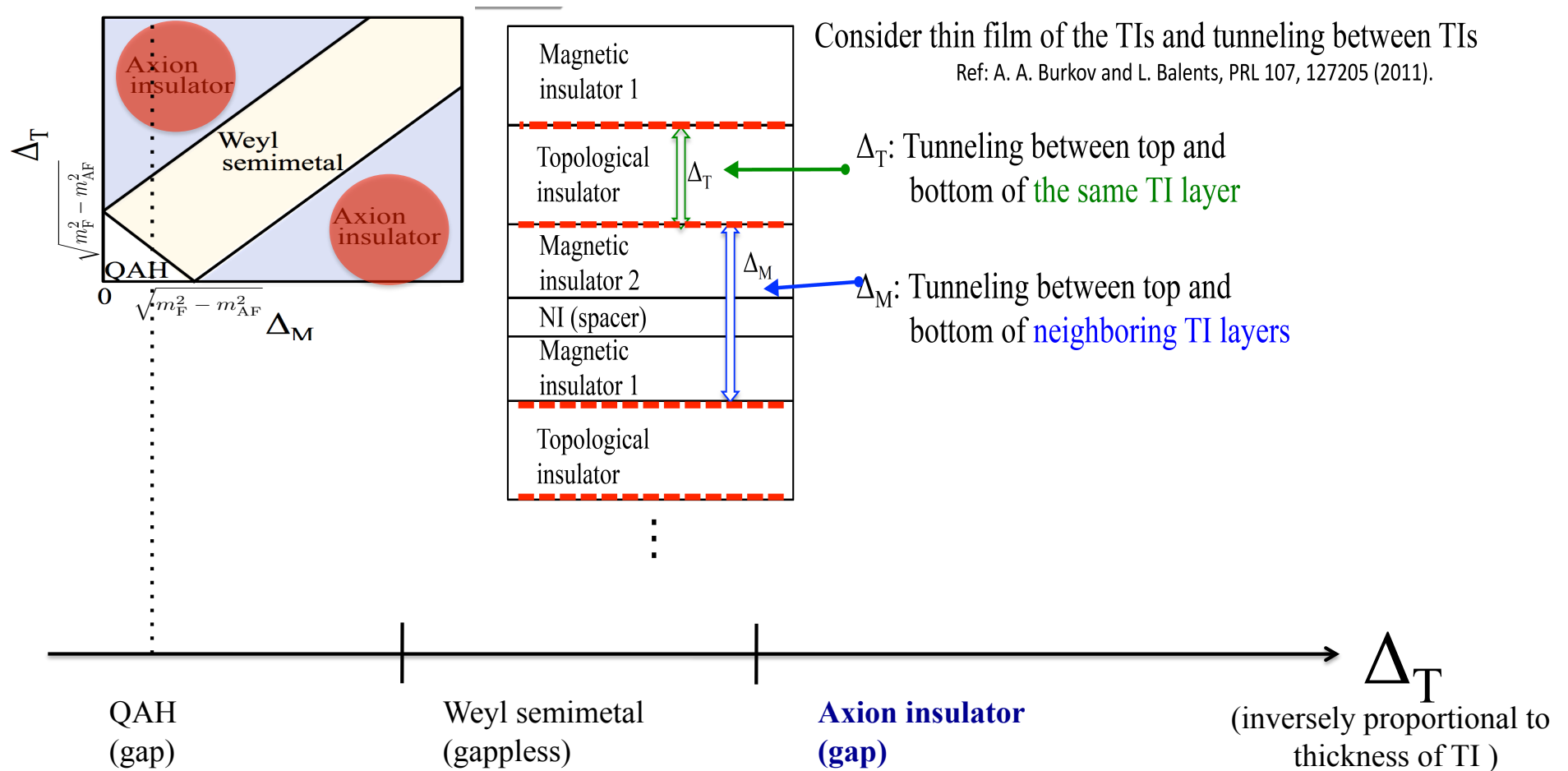
Model

To control $\theta = \text{Arctan}(\varphi_5 / \varphi_4)$

we consider the superlattice of topological insulators (TIs) and magnetic insulators (for P and T symmetry breaking)

Superlattice of magnetic-topological insulator (TI) and magnetic insulator (MI1 and 2)





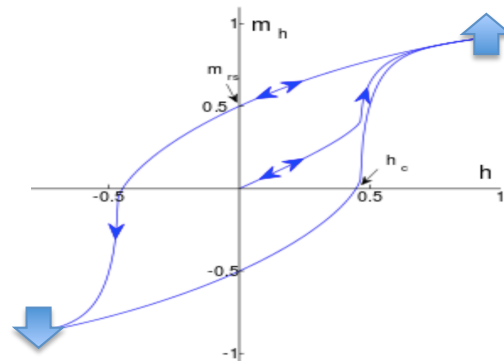
We consider axion insulator

How to derive $\partial_t \theta \neq 0$

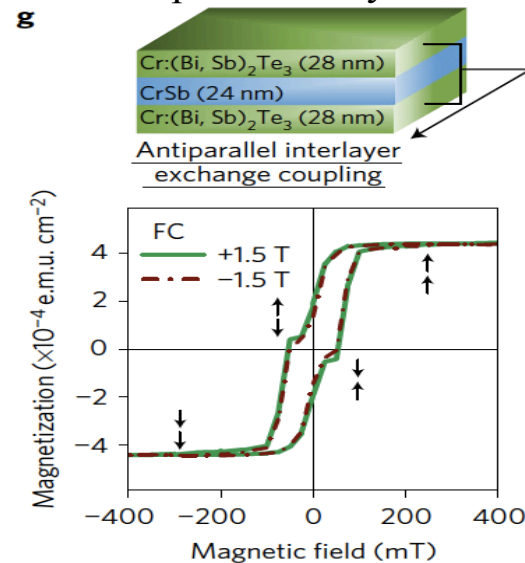
Anti-magnetic configuration is important for manipulation of θ

The magnetic order can be manipulated by an applied magnetic field B_{ex}

“Hard” M-H loop



M-H loop in multilayer



Q. L. He et. al., Nat Mater 16, 94 (2017)

Spin control of super-lattice device

Assume:

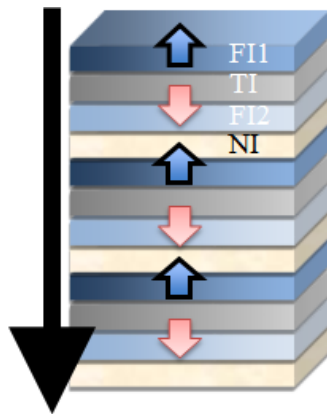
M1 soft

$M1 \neq M2$

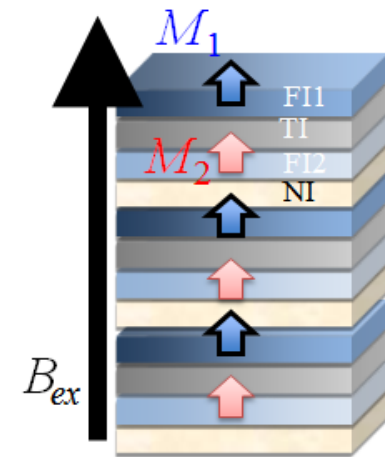
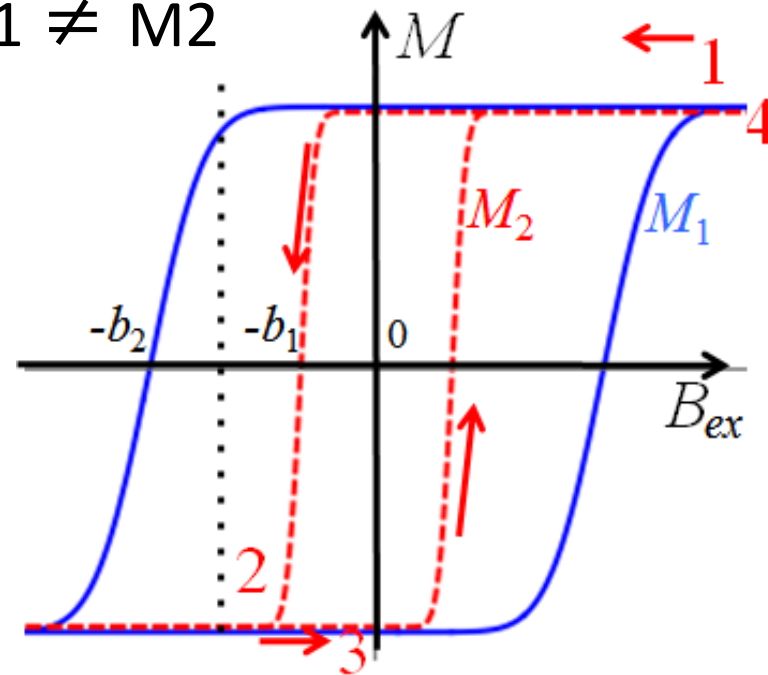
M2 hard

$$\theta = \text{Arctan}(\varphi_5 / \varphi_4)$$

$$b_1 < B_{ex} < b_2$$



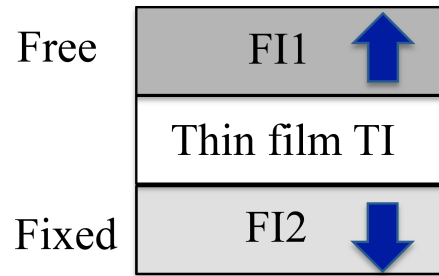
⇒ starting point



Simple form

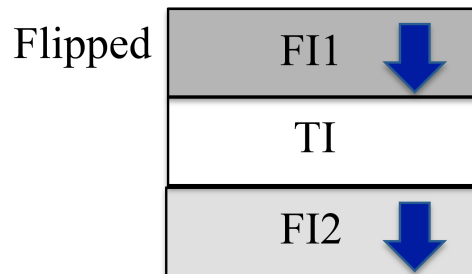
$$\theta = 0$$

In the absence of B_{ex} : $\delta\theta \neq 0$



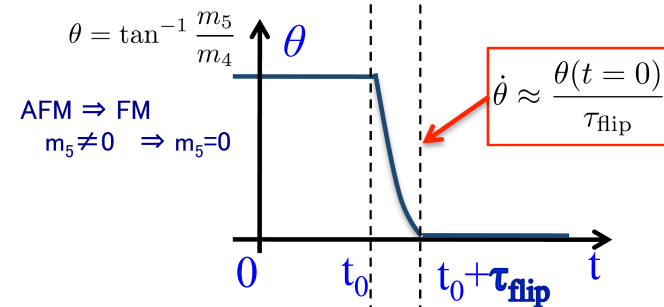
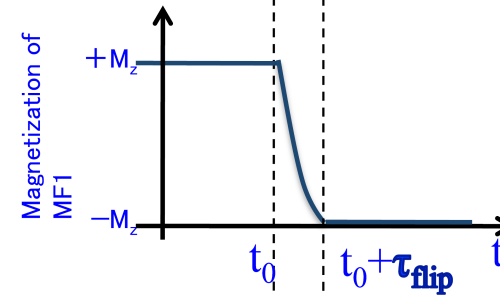
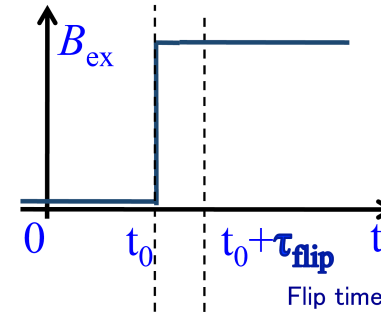
$t < t_0$ (initial state)

$B_{\text{ex}} = 0$



$t > t_0 + \tau_{\text{flip}}$

Strong B



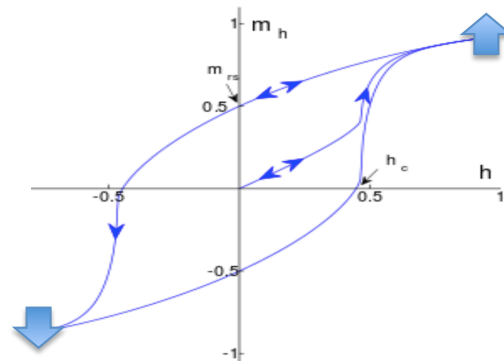
$$m_{\text{AF}} = m_{\text{AF}}(t=0) \exp\left(-\frac{t}{\tau_f}\right)$$

How to derive $\partial_t \theta \neq 0$

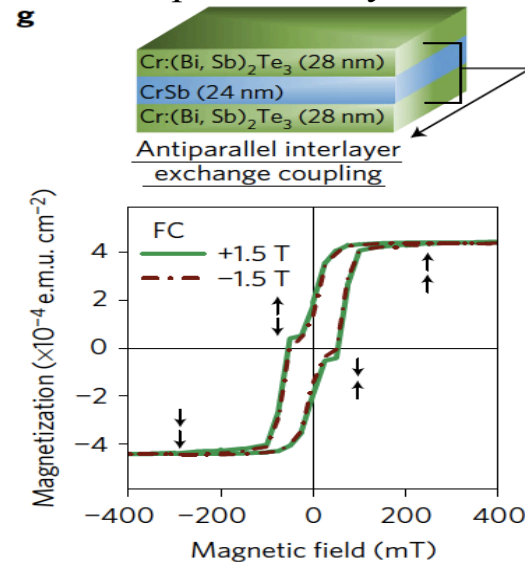
Anti-magnetic configuration is important for manipulation of θ

The magnetic order can be manipulated by an applied magnetic field B_{ex}

“Hard” M-H loop



M-H loop in multilayer




Q. L. He et. al., Nat Mater 16, 94 (2017)

Derivation

We assume the following Lagrangian

$$\begin{aligned}\mathcal{L} &= \mathcal{L}_{\text{Maxwell}} + \mathcal{L}_{\text{Axion}} + \mathcal{L}_e \\ &= \frac{1}{2} \left(\epsilon E^2 - \frac{1}{\mu} B^2 \right) + \frac{\alpha}{2\pi\hbar} \theta \mathbf{E} \cdot \mathbf{B} + \mathbf{j}_e \cdot \mathbf{A}\end{aligned}$$



Assume

$$\begin{aligned}\mathbf{B} &= \mu_0(\mathbf{H} + \mathbf{M}) \rightarrow \mu_0 \mathbf{H} \\ \mathbf{D} &= \tilde{\epsilon} \epsilon_0 \mathbf{E} \rightarrow \epsilon_0 \mathbf{E}\end{aligned}$$



Maxwell equations

$$\begin{aligned}\nabla \times \mathbf{E} &= -\partial_t \mathbf{B} \\ \nabla \times \mathbf{H} &= \partial_t \mathbf{D} + \mathbf{j}_e + \mathbf{j}_a \\ \mathbf{j}_a &= -\frac{c\epsilon_0\alpha}{\pi} [(\nabla\theta) \times \mathbf{E} + (\partial_t\theta)\mathbf{B}]\end{aligned}$$

α : fine structure, c : velocity of light

We obtain the wave equation for \mathbf{B}

$$\left(\frac{\partial^2}{\partial t^2} - c^2 \nabla^2 \right) \mathbf{B}_{\text{ind}} = \frac{1}{\epsilon_0} \nabla \times \mathbf{j}$$

↑ Induced magnetic field
 ↗ Source term

$$\frac{1}{\epsilon} \nabla \times \mathbf{j} = \frac{1}{\epsilon} \nabla \times \mathbf{j}_e - \frac{c\alpha}{\pi} [(\nabla \times \mathbf{E}) \nabla \theta - (\nabla \theta) \cdot \nabla \mathbf{E}] - \frac{c\alpha\mu}{\pi} \partial_t \theta (\partial_t \mathbf{D} + \mathbf{j}_e + \mathbf{j}_a), \quad (1)$$

Assume

$$\nabla \times \mathbf{j}_e = 0 \text{ and } \mathbf{E} \text{ is spatial uniform}$$

Result:

$$\left(\frac{\partial^2}{\partial t^2} - c^2 \nabla^2 \right) \mathbf{B}_{\text{ind}} = -\frac{c\alpha\mu_0}{\pi} \partial_t \theta (\partial_t \mathbf{D} + \mathbf{j}_e + \mathbf{j}_a)$$

$$\mathbf{j}_a = -\frac{c\epsilon_0\alpha}{\pi} [(\nabla \theta) \times \mathbf{E} + (\partial_t \theta) \mathbf{B}]$$

Main message:

dynamical θ is driven by applied magnetic field (B_{ex})

Result:

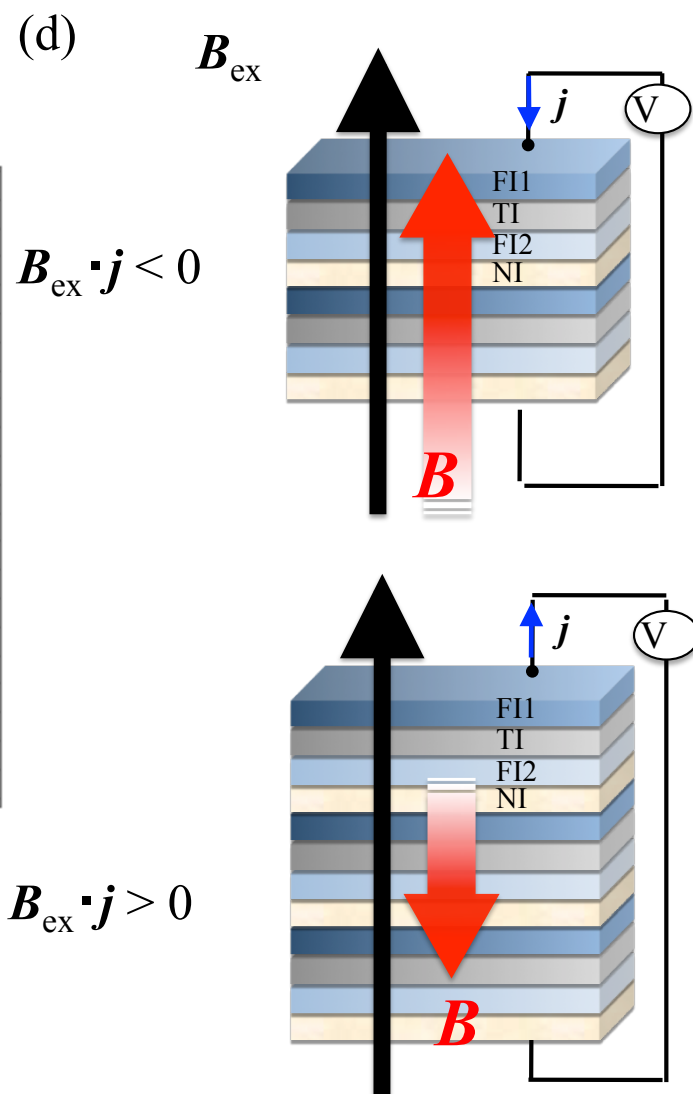
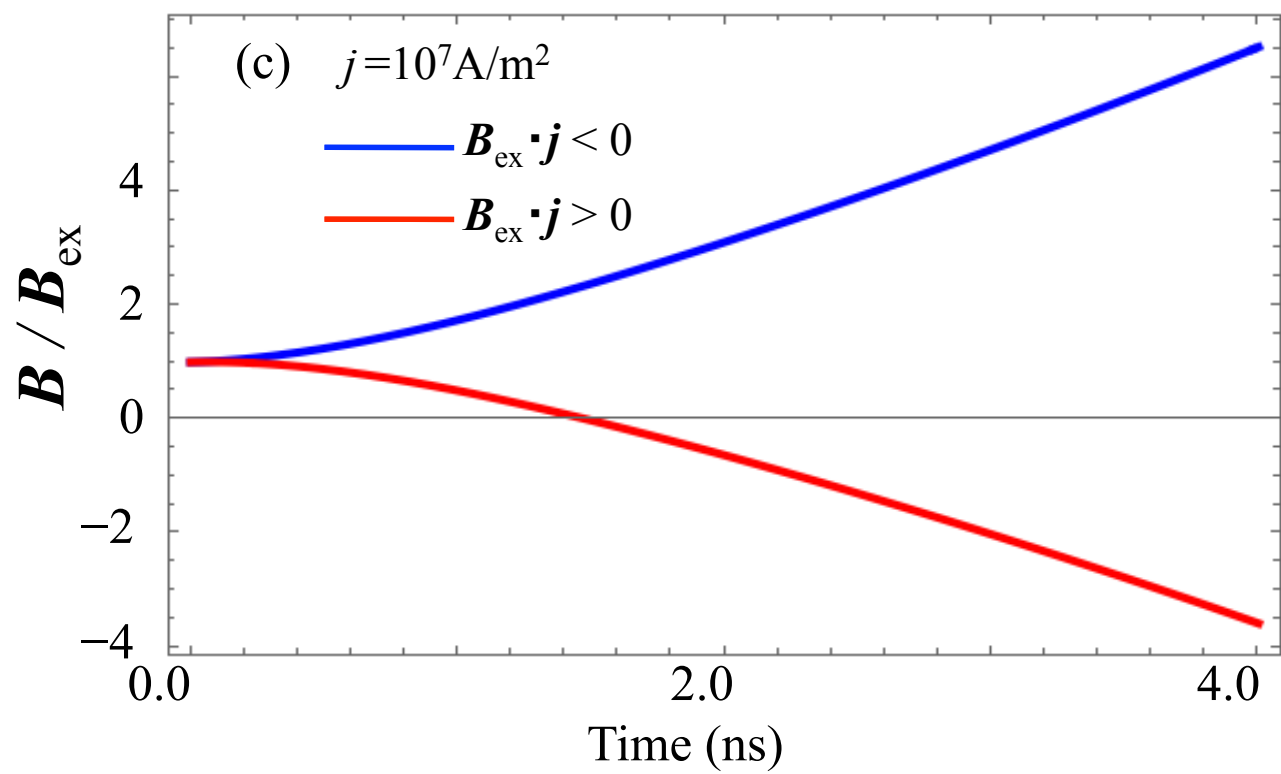
$$\left(\frac{\partial^2}{\partial t^2} - c^2 \nabla^2 \right) \mathbf{B}_{\text{ind}} = - \frac{c\alpha\mu_0}{\pi} \partial_t \theta (\partial_t \mathbf{D} + \mathbf{j}_e + \mathbf{j}_a)$$

When we apply the external magnetic field,
the dynamical θ is induced
and
the dynamical θ drives the electromagnetic effects

(Case 1) in the presence of only the external magnetic field

(Case 2) in the presence of the external magnetic field and charge current (*layered direction*)

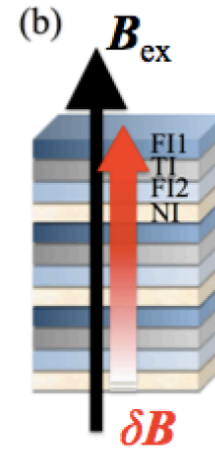
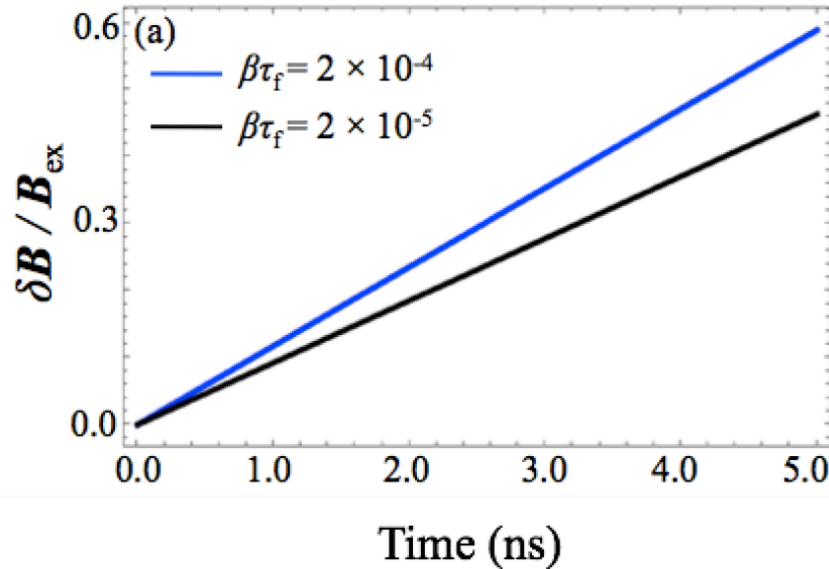
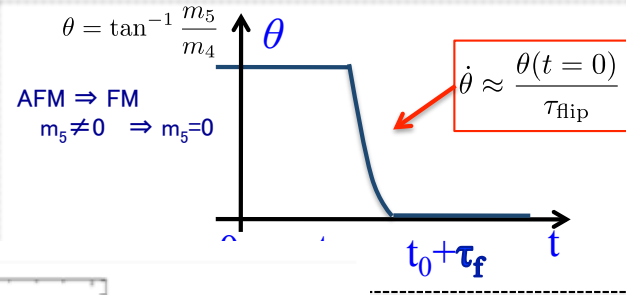
(Case 3) in the presence of the external magnetic field and charge current (*inplane*)



(Case 1) in the presence of only the external magnetic field B_{ex}

$$\left(\frac{\partial^2}{\partial t^2} - c^2 \nabla^2\right) \mathbf{B}_{\text{ind}} = -\frac{c\alpha\mu_0}{\pi} \partial_t \theta (\partial_t \mathbf{D} + \mathbf{j}_e + \mathbf{j}_a)$$

$$= \frac{\alpha^2}{\pi^2} (\partial_t \theta)^2 \mathbf{B}_{\text{ex}}$$

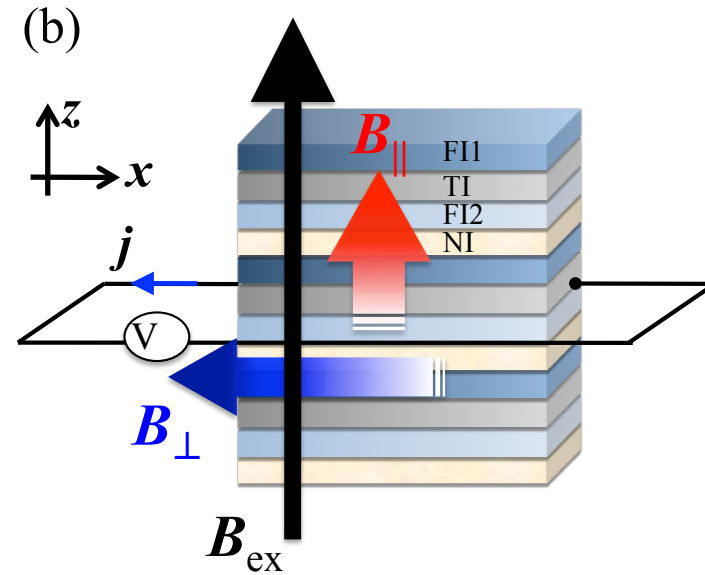
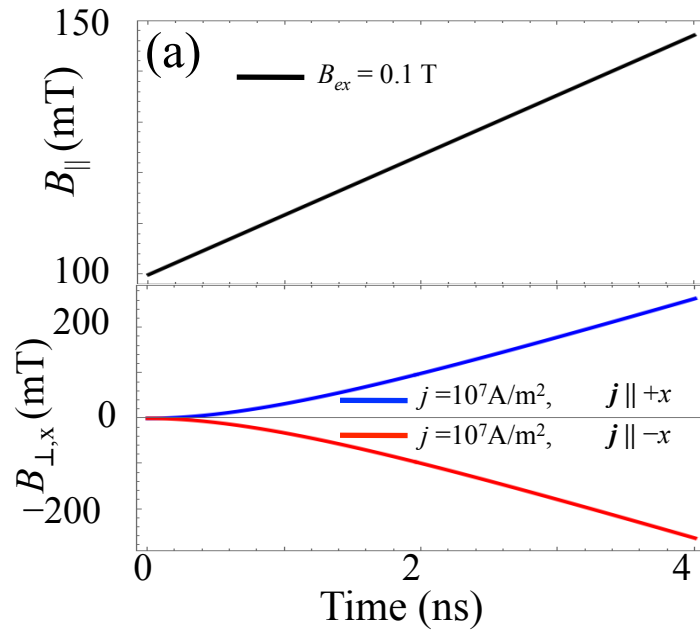


Additionally induced
magnetic field δB

Besides, charge current
via *chiral magnetic effect*

(Case 3) in the presence of the external magnetic field and charge current (*inplane*)

$$\left(\frac{\partial^2}{\partial t^2} - c^2 \nabla^2 \right) \mathbf{B}_{\text{ind}} = \underbrace{-\frac{c\alpha\mu_0}{\pi} \partial_t \theta (\partial_t \mathbf{D} + \mathbf{j}_e)}_{\substack{\text{Contribution} \\ \text{from} \\ \text{charge current}}} + \underbrace{\frac{\alpha^2}{\pi^2} (\partial_t \theta)^2 \mathbf{B}_{\text{ex}}}_{\substack{\text{Contribution} \\ \text{from } \mathbf{B}_{\text{ex}}}}$$



How can this effect be detected ?

[1]. Planar hall effect

Planar hall voltage $\propto M^2 \sin\theta$

\Rightarrow provably μV order and it's possible to detect the nsec response

[2]. Chiral magnetic effect

$$\text{rot}B - \frac{1}{c} \frac{\partial E}{\partial t} = 4\pi \left[j + \frac{e^2}{2h} \nabla \left(\frac{\theta}{\pi} \right) \times E + \frac{e^2}{2\pi h} \frac{\dot{\theta}}{\pi} B \right]$$



$$j \propto \dot{\theta} B \quad \text{Chiral magnetic effect}$$

Probably, device resistance should be $\sim \text{M}\Omega$, it's possible to detect $< \text{nA}$ current.

Summary

[1] We study the multilayer of topological insulators (TI) and magnetic insulator (MI)

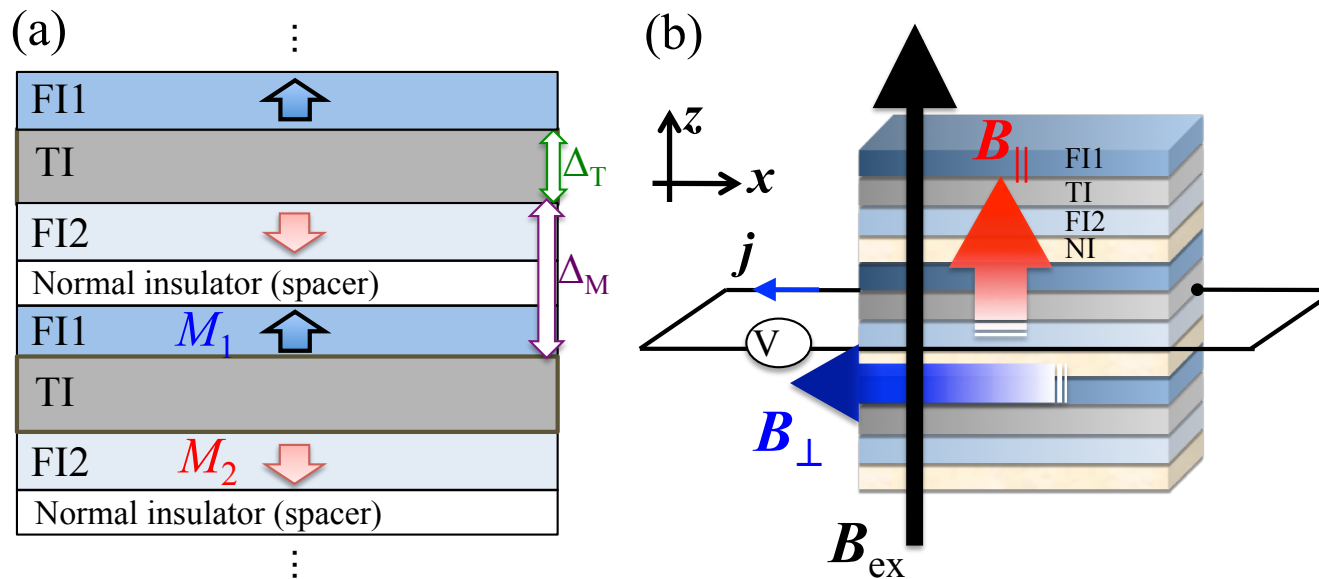
$$\theta = \text{Arctan}(\varphi_5 / \varphi_4)$$

[2] For “axion”-response in materials,

we consider axion insulator and how to drive $\partial_t \theta \neq 0$

⇒ Use magnetization flip by B_{ex}

[3] **Main result** We find that $\partial_t \theta$ and an applied electric field trigger magnetic field [Electromagnetic effect]



Summary

- Axion is also important in the material science, specially in topological insulator $\theta = \text{Arctan}(\varphi_5 / \varphi_4)$
- By this collaboration, we phenomenologically found new method and response for dynamical axion response
- Next plan is to first detect this response and connect this collaboration to particle physics and cosmology

Please join us if you are interested in such activities !!